# **ELECTROMAGNETISM SUMMARY**

- ► Electrostatics
- Charge conservation
- > Dielectrics
- Electromagnetic Shielding

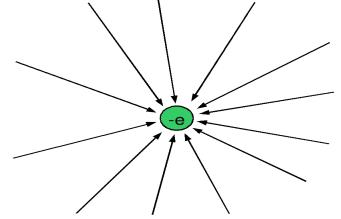


## **Electrostatics: The Case of Stationary Charge**

- > The source of all electromagnetic fields is ultimately the charge.
- > When there are no time variations, charge is the source of electric field.
- ► For a point charge we have Coulomb's law:

$$E_r = \frac{q}{4\pi\epsilon_0 r^2}$$

where the free space permittivity,  $\epsilon_0 = 8.85 \times 10^{-12}$  Farads/m and q is in *Coulombs*.





# **Electrostatics: The static electric field**

In general we have the following closed integral form (Gauss's law):  $\oint \mathbf{E}.\mathbf{dA} = \frac{q}{\epsilon_0}$ > Example: a charged capacitor.  $E = \sigma / \epsilon_o$ ... Area lack of electrons E field net electrons  $\varepsilon$  E x Area =  $\sigma$  x Area



## **Electrostatics: The Electrostatic Potential**

Definition: The potential difference between two points x<sub>1</sub> and x<sub>2</sub> is given by,

$$\Phi = -\int\limits_{\mathbf{x}_1}^{\mathbf{x}_2} \mathbf{E}.\mathbf{dl} = -\int\limits_{\gamma} \mathbf{E}.\mathbf{dl}$$

- Since the path  $\gamma$  can be any which connects the points  $x_1$  and  $x_2$  we may conclude that  $\mathbf{E} = -\nabla \Phi$ .
- Kirchhoffs voltage law.



## **Charge Conservation**

When current flows out of a region in space, it depletes the charge in that region. The current per unit area j is given by,

$$I = \oint_{A} \mathbf{j}.\mathbf{dA} = -\frac{\partial q}{\partial t}$$

where the current I is,  $I = \int \mathbf{j} \cdot \mathbf{dA}$  for any surface (not necessarily closed).

- If q is in Coulombs and the time in seconds then I is in Amperes.
- ► A general law... not just electrostatics.
- If current flows round in a closed loop, then there need be no change in the charge: Kirchhoffs current law.
- Quite generally: j = nqv where n is the charge carrier density, q their charge and v their velocity.



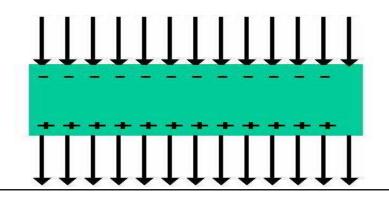
#### **Dielectrics and Conductors**

- Dielectrics are insulating materials that do not allow D.C. current to flow through them. Usually we just call them insulators.
- ▶ In conductors, charge carriers (electrons) are free to move. E.G. metals.
- We study briefly the phenomenology of dielectrics and conductors.



## **Dielectrics 1**

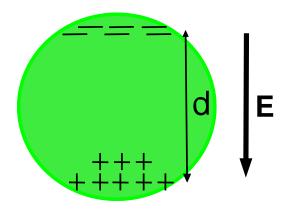
- Dielectrics are insulating materials that do not allow D.C. current to flow through them.
- Electrons and nuclei in the atoms and molecules of dielectrics experience opposing forces in the presence of an imposed electric field.
- Electrons move opposite to the field and nuclei move in the direction of the field. This separation of charge produces a **polarisation**.
- The charge separation induced by the field, acts to reduce the electric field within the dielectric.





#### **Dielectrics 2: Polarisation**

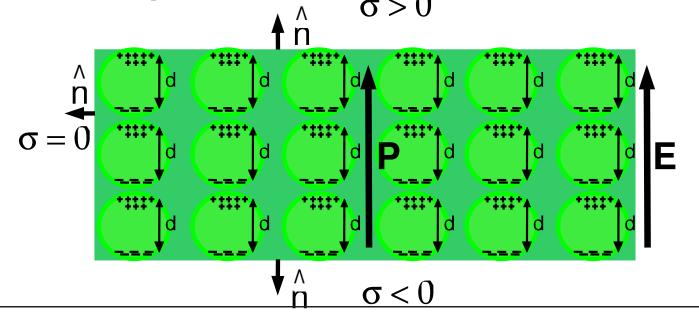
- Polarisation P is the dipole moment per unit volume induced by an imposed, external electric field
- > P = nqd, where n is the number density of dipoles, q is the charge at each end of the dipole and d is the displacement of the  $\pm q$  charges at each end of the dipole.
- Polarisation is a vector quantity.
- > Pictorial representation of a dipole moment:





#### **Dielectrics 3: Polarisation**

- At the edge of a polarised dielectric there is a charge density left over by the displacement of the dipoles.
- > The surface charge density on the dielectric  $\sigma_p = \mathbf{P}.\hat{\mathbf{n}}$  where  $\hat{\mathbf{n}}$  is unit vector normal to the surface.
- > Notice that  $\sigma_p$  belongs to the material. It is not a **free** charge.





## **Dielectrics 4: Relative dielectric constant**

> The relative dielectric constant  $\epsilon_r$  is defined by,

 $\mathbf{P} = \epsilon_0(\epsilon_r - 1)\mathbf{E}$ 

► The Electric Displacement D is defined by,

 $\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P} = \epsilon_r \epsilon_0 \mathbf{E}$ 

Main advantage of the definition of D is that its source is the free charge only and not the induced polarisation charge,

$$\oint_{A} \mathbf{D}.\mathbf{dA} = q_{free}$$

$$\blacktriangleright \text{ c.f. } \oint_{A} \mathbf{E}.\mathbf{dA} = \frac{q_{total}}{\epsilon_0}$$



#### **Dielectrics 5: Key Points**

- > The relative dielectric constant  $\epsilon_r$  describes the behaviour of a dielectric when exposed to an oscillating electric field.
- >  $\epsilon_r$  at D.C. is a positive dimensionless number and  $\epsilon_r > 1$  for dielectrics.
- > When the electric field oscillates,  $\epsilon_r$  is a complex function of frequency.
- > The ratio of the imaginary to real components of  $\epsilon_r$  is termed the **loss** tangent of the dielectric:

$$tan\delta = \frac{Im(\epsilon_r)}{Re(\epsilon_r)}$$



## **Conductors and Ohm's law**

Ohm's law: The current density in a conductor is proportional to the electric field within the conductor.

 $\mathbf{j} = \sigma \mathbf{E}$ 

where  $\sigma$  is the conductivity.

- > Conductors are completely specified by  $\sigma$ .
- > For copper,  $\sigma = 5.80 \times 10^7 mhos/meter$ ..
- Ohm's law is assumed to be an accurate result for metals at all radiofrequencies :).
- > I.E.  $\sigma$  is always a real number and independent of frequency.



## **Conductors vs Dielectrics?**

- > For metals:  $\mathbf{j} = \sigma \mathbf{E}$
- > For dielectrics:  $P = \epsilon_0(\epsilon_r 1)E$
- ► If P oscillates as a function of time then,

 $j\omega \mathbf{P} = j\omega\epsilon_0(\epsilon_r - 1)\mathbf{E}$ 

- Since the dipoles are reversing sign at rate  $\omega$  while traversing a distance d we may write,  $\mathbf{j}_{\mathbf{P}} = \mathbf{j}\omega\mathbf{P}$ , where  $\mathbf{j}_{\mathbf{P}}$  is the polarisation current.
- > Thus dielectrics obey a sort of Ohm's law with  $\mathbf{j}_{\mathbf{P}} = \sigma_{\mathbf{P}} \mathbf{E}$  and  $\sigma_{P} = j\omega\epsilon_{0}(\epsilon_{r} 1)$
- > The main difference between conductors and insulators is simply that  $\sigma$  is **resistive** for a conductor, but is mainly **reactive** for insulators.
- > In fact, generally speaking,  $\epsilon_r$  is a complex function of frequency.



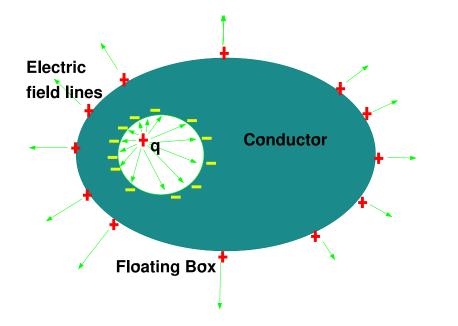
**Electromagnetic Shielding 1** 

Ohm's law for metals and Gauss's law for the electric field give rise to the concept of electromagnetic shielding



# **Electromagnetic Shielding 2: The Floating Conducting Box**

Suppose that there is an isolated positive charge, q, in a cavity and that the conducting box is **floating** (not connected to earth) and **neutral** (has no net charge).





# **Electromagnetic Shielding 3: The Floating Conducting Box**

We can deduce the following...

- Electric current is finite (actually zero here) and Ohm's law implies that the electric field inside the conductor (Turquoise region) is zero.
- Gauss's law implies that the net negative charge on the inside of the inner wall of the cavity is equal and opposite in sign to q.
- Charge conservation implies that the amount of positive charge on the outer surface equals the negative charge on the inner surface.



# **Electromagnetic Shielding 4: The Floating Conducting Box**

Observations:

- (a) The negative charge arranges itself to give perfect cancellation of the electric field inside the conductor.
- (b) The negative surface charge density ( $\sigma$ ) is highest on the wall closest to q.
- (c) The positive charges on the outer surface see no electric field inside the conductor and therefore do not respond to movement in either q or the negative charge on the cavity wall.
- (d) The positive charges on the outside bunch toward surfaces of high curvature and spread out along surfaces of low curvature.



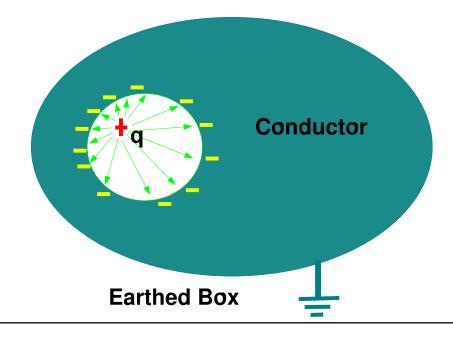
## **Electromagnetic Shielding 5: The Floating Conducting Box**

- (e) The charge on the surface of the metal-air interface lies right on the surface and not inside the conductor. (RF burns?)
- (f) By Gauus's law, here is still an electric field outside the conductor.
- (g) If there were no net charge inside the cavity (q = 0) then there would be **no electric field outside the conducting box either**. For example there could even be a +q and a -q arbitrarily located. It makes no difference.
- (h) At the surface of the metal, the electric field vector is always normal. There is **never** any tangential component of electric field on the surface of a conductor.



# **Electromagnetic Shielding 6: The Earthed Conducting Box**

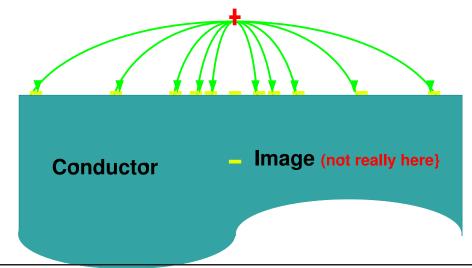
- Largely the same observations as for the floating case, except that now there is **no field outside the conductor** under any condition.
- Main practical conclusion: Closed conducting boxes isolate the outside world from electromagnetic influences within the box and vice versa.





# **Electromagnetic Shielding 7: Proximity to a Conductor**

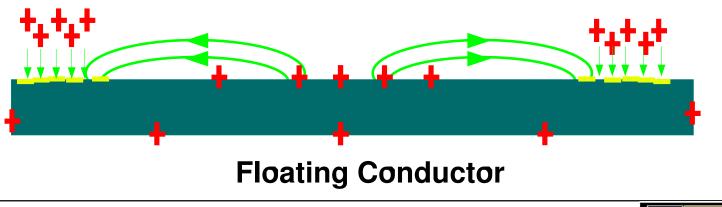
- A charge near a conducting surface attracts charge of the opposite sign to the nearest point on the surface.
- These surface charges arrange themselves so that the tangntial component of electric field is zero on the surface.
- One may compute the electric field outside the conductor by assuming (mathematically) that there is an **image** charge within the conductor.





# **Electromagnetic Shielding 8: Proximity to Floating Conductors**

- Two charges (electronic components) near a floating conductor must share the charge within the conductor.
- Electromagnetic fields at opposite ends of the plane cause charge separations that increase the influence of each component on the other.
- If the electronic components are themselves neutral then there is not such a large influence of the plane.





## **Electromagnetic Shielding 9: Proximity to Earthed Conductors**

- > Earth is an inexhaustible source of charge.
- Two charges (electronic components) near an earthed conductor attract charge from earth.
- Electromagnetic fields at opposite ends of the plane cause charge separations that **decrease** the influence of each component on the other.

